

Curl in Coordinate Systems

Consider now the curl of vector fields expressed using our coordinate systems.

Cartesian

$$\begin{aligned}\nabla \times \mathbf{A}(\bar{r}) = & \left[\frac{\partial A_y(\bar{r})}{\partial z} - \frac{\partial A_z(\bar{r})}{\partial y} \right] \hat{a}_x \\ & + \left[\frac{\partial A_z(\bar{r})}{\partial x} - \frac{\partial A_x(\bar{r})}{\partial z} \right] \hat{a}_y \\ & + \left[\frac{\partial A_x(\bar{r})}{\partial y} - \frac{\partial A_y(\bar{r})}{\partial x} \right] \hat{a}_z\end{aligned}$$

Cylindrical

$$\begin{aligned}\nabla \times \mathbf{A}(\bar{r}) = & \left[\frac{1}{\rho} \frac{\partial A_z(\bar{r})}{\partial \phi} - \frac{\partial A_\phi(\bar{r})}{\partial z} \right] \hat{a}_\rho \\ & + \left[\frac{\partial A_\rho(\bar{r})}{\partial z} - \frac{\partial A_z(\bar{r})}{\partial \rho} \right] \hat{a}_\phi \\ & + \left[\frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho A_\phi(\bar{r})) - \frac{1}{\rho} \frac{\partial A_\rho(\bar{r})}{\partial \phi} \right] \hat{a}_z\end{aligned}$$

Spherical

$$\begin{aligned} \nabla \times \mathbf{A}(\bar{r}) = & \left[\frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\phi(\bar{r})) - \frac{1}{r \sin \theta} \frac{\partial A_\theta(\bar{r})}{\partial \phi} \right] \hat{a}_r \\ & + \left[\frac{1}{r \sin \theta} \frac{\partial A_r(\bar{r})}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi(\bar{r})) \right] \hat{a}_\theta \\ & + \left[\frac{1}{r} \frac{\partial}{\partial r} (r A_\theta(\bar{r})) - \frac{1}{r} \frac{\partial A_r(\bar{r})}{\partial \theta} \right] \hat{a}_\phi \end{aligned}$$

Yikes! These expressions are **very** complex. Precision, organization, and patience are required to **correctly** evaluate the **curl** of a vector field!